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In this note we shortly report on the results of a calculation of the spectra and probabilities of $\eta^0 \rightarrow \gamma + e^+ + e^-$ and $\eta^0 \rightarrow \gamma + \mu^+ + \mu^-$. We assume for η^0 the quantum numbers $J = 0$, $P = -1$, $I = 0$, $C = G = +1$. We approximate the $\eta^0 \rightarrow 2\gamma$ amplitude (with one photon off-mass-shell) with a subtraction term and the pole terms due to ρ^0 and ω^0 :

(1)

We call m the invariant mass of the emitted two-lepton system

(2)
$$m = \left[(E_+ + E_-)^2 - (p_+ + p_-)^2 \right]^{1/2}$$

where E_{\pm} and p_{\pm} are the energies and momenta of the positive (negative) lepton. The possible values of m are between $2m_l$ (where m_l is the lepton mass) and m_{η} (η - mass $\simeq 550$ MeV). The number of events with m between m and $m + dm$ is given by:

(x) - Istituto di Fisica dell'Università di Firenze.

$$\begin{aligned}
 N(m) dm = & \left[\tau(2\gamma) \right]^{-1} \frac{4\alpha}{3\pi} \frac{1}{m} \left[1 - \left(\frac{m}{m_\eta} \right)^2 \right]^3 \left[1 + 2 \left(\frac{m}{m} \right)^2 \right] \times \\
 (3) & \times \left[1 - \left(\frac{2m}{m} \right)^2 \right]^{1/2} \left[v \frac{m_v^2}{m_v^2 - m^2} + (1 - v) \right]^2 dm
 \end{aligned}$$

where v is a parameter depending on the relative values of the residues of the poles and the subtraction constant. In (3) $\tau(2\gamma)$ is the partial lifetime for $\eta^0 \rightarrow 2\gamma$. The mass m_v is some average of the ρ^0 and ω^0 mass. The value $v=0$ corresponds to keeping only the subtraction term in the expansion (1) (constant form factor). The value $v=1$ corresponds to keeping only the vector meson poles neglecting the subtraction constant. There exists an experiment on the π^0 form-factor that gives the value of the derivative at the origin of the π^0 form-factor with respect to the squared four-momentum of the off-mass-shell photon⁽¹⁾. If we assume that η^0 is the eighth member of the unitary symmetry octet containing π and K we can tentatively make use of unitary symmetry to obtain a value for v from the quantity measured in the π^0 experiment. We get $v = -7 \pm 5$. Needless to say, such an extrapolation of unitary symmetry arguments to a low energy region where mass differences are quite important may be completely illusory. In fig. 1 we have reported the entire spectrum of $\eta^0 \rightarrow \gamma + \mu^+ + \mu^-$ and the high energy tail of $\eta^0 \rightarrow \gamma + e^+ + e^-$ in arbitrary units, for values of $v = 0, 1$ and -7 . The spectra have the right relative normalizations, i.e., apart from the common factor $(4\alpha/3\pi)(\tau(2\gamma))^{-1}$, $N(m)dm$ gives, for each case, directly the number of events with m between m and $m + dm$.

Of course, for $\eta^0 \rightarrow \gamma + e^+ + e^-$ the spectrum receives its biggest contribution from smaller values of m down to $2m_e$, as given by (3). The branching ratios ρ_e and ρ_μ for $\eta^0 \rightarrow \gamma + e^+ + e^-$ relative to $\eta^0 \rightarrow 2\gamma$ can be obtained by integrating (3). For $v=0$ one has

$$\begin{aligned}
 \rho_\mu = & \frac{2\alpha}{3\pi} \left\{ \left(-\frac{7}{2} + 13r^2 + 4r^4 \right) (1 - 4r^2)^{1/2} + \right. \\
 (4) & \left. + 2(1 - 18r^4 + 8r^6) \lg \frac{1 + (1 - 4r^2)^{1/2}}{2r} \right\}
 \end{aligned}$$

with $r = m_\mu/m_\eta$. Putting $r=0$ one obtains

$$\rho_e = \frac{2\alpha}{3\pi} \left[\log \left(\frac{m}{m_e} \right)^2 - \frac{7}{2} \right],$$

(1) - N. P. Samios - Phys. Rev. 121, 275 (1961).

which is the well-known Dalitz formula.

By numerical integration, with $m_V = 750$ MeV, we find

$$(5) \quad \rho_e = (16.2 + 0.47 v + 0.035 v^2) \times 10^{-3}$$

$$(5') \quad \rho_\mu = (55.8 + 21.9 v + 2.74 v^2) \times 10^{-5}$$

From (5) and (5') we see that ρ_e is much less sensitive to v than ρ_μ and, taking the model literally, we expect, independently of v , $\rho_e > 14.6 \times 10^{-3}$ and $\rho_\mu > 12.1 \times 10^{-5}$. With $v = 0$ $\rho_e = 16.2 \times 10^{-3}$ and $\rho_\mu = 55.8 \times 10^{-5}$. With $v = 1$ $\rho_e = 16.7 \times 10^{-3}$ and $\rho_\mu = 80.4 \times 10^{-5}$. With $v = -1$ (unitary symmetry extrapolation) $\rho_e = 14.6 \times 10^{-3}$ and $\rho_\mu = 36.9 \times 10^{-5}$. Note that ρ_e alone would determine v (or better two possible values for v) and then one would be able to predict ρ_μ and the shapes of the spectra. A reaction such as $\gamma + p \rightarrow \eta^0 + p$ followed by $\eta^0 \rightarrow \mu^+ + \mu^- + \gamma$ simulates $\gamma + p \rightarrow \mu^+ + \mu^- + \gamma + p$. With a cross-section $\sim 10^{-30}$ cm² for $\gamma + p \rightarrow \eta^0 + p$ the apparent cross-section for $\gamma + p \rightarrow \mu^+ + \mu^- + \gamma + p$ would then be at least $\sim 10^{-34}$ cm².

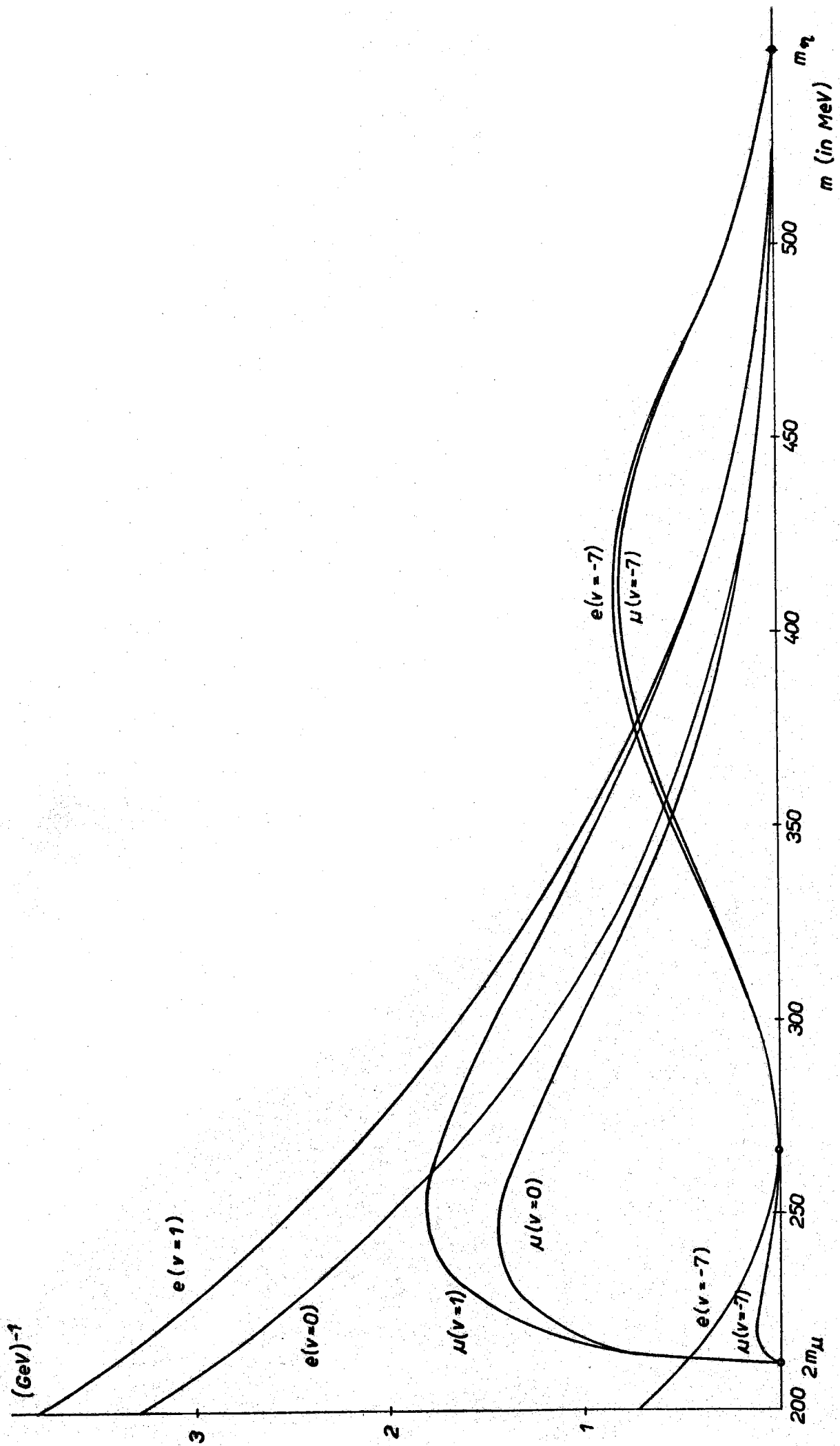


FIG. 1 - Graphs of $\frac{3\pi}{4d} \mathcal{Z}(2\gamma) N(m)$, where $N(m)dm$ is the number of events with m (invariant lepton mass) between m and $m+dm$, for $\eta^0 \rightarrow \gamma + \mu^+ + e^-$ and $\eta^0 \rightarrow \gamma + \mu^+ + \mu^-$ with different values of the parameter ν .